

SHORT CONTRIBUTION

On the evaluation of the wet bulb temperature as a function  
of dry bulb temperature and relative humidity

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(Manuscript received June 18, 1993; accepted in final form January 17, 1994)

RESUMEN

Se propone un polinomio de grado tres para evaluar la presión de vapor de saturación ( $e_s$ ), a fin de convertir la ecuación psicrométrica en una ecuación cúbica en  $T_w$  (temperatura de bulbo húmedo), puesto que no es posible expresar a  $T_w$  como una función directa de la temperatura de bulbo seco ( $T$ ), la humedad relativa ( $RH$ ) y la presión ( $p$ ). Dicha expresión para  $e_s$  es más precisa que otras semiempíricas reportadas en la literatura. Los resultados del cálculo de  $T_w$  mediante la ecuación cúbica se compararon con los valores arrojados por un método iterativo, resultando ser satisfactorios.

ABSTRACT

This paper proposes the use of a cubic polynomial to evaluate the saturation vapor pressure ( $e_s$ ) in order to transform the psychrometric equation into a cubic equation in  $T_w$  (wet bulb temperature). This is done because it is not possible to write  $T_w$  as a direct function of dry bulb temperature ( $T$ ), relative humidity ( $RH$ ) and pressure ( $p$ ). Statistical comparisons show that this polynomial is accurasier than other semi-epirical methods used to calculate  $e_s$ . Results of the third degree equation for  $T_w$  were compared whit resultant values of an iterative method and they were satisfactory.

1. Introduction

The Clausius-Clapeyron equation establishes a relationship between the saturation vapor pressure ( $e_s$ ) and the dry bulb temperature ( $T$ ). If the process involves only the vapor and water phases, the Clausius-Clapayron expression for temperatures between  $-10^{\circ}\text{C}$  to  $50^{\circ}\text{C}$  is (Bindon, 1965):

$$e_s = \exp\left(\frac{21.4T + 494.41}{T + 273.15}\right) \quad (1)$$

where the  $e_s$  units are hectopascals (hPa).

It is possible to express the relative humidity ( $RH$  in %) in terms of the atmospheric pressure ( $p$  in hPa), the dry bulb temperature ( $T$  in  $^{\circ}\text{C}$ ), and the wet bulb temperature ( $T_w$  in  $^{\circ}\text{C}$ ) from the psychrometric relation:

$$e_{sw} - \frac{RH}{100}e_s = kp(T - T_w) \quad (2)$$

where  $e_{sw}$  is Eq. 1 evaluated in terms of  $T_w$  instead of  $T$ , and  $K$  is  $1/0.622$  times the ratio of the water vapor specific heat at constant pressure to the latent heat of condensation. The value of  $k$  however depends on the psychrometer type. Bindon (1965) has reviewed the  $k$  values of different psychrometer types obtained by various authors. A summary is presented in Table 1. The theoretical value of  $6.53 \times 10^{-4} \text{ }^{\circ}\text{C}^{-1}$  will be used in this paper for evaluation of psychrometric variables.

Table 1. Psychrometric constants for water ( $K_w$ ) and ice ( $K_i$ ) from some authors reported by Bindon (1965).

$K_w$ ( $^{\circ}\text{C}^{-1}$ )	$K_i$ ( $^{\circ}\text{C}^{-1}$ )	Ventilation	Table or slide rules
$6.53 \times 10^{-4}$	$5.703 \times 10^{-4}$	Almost 3 m/sec	Thermodynamic
$6.62 \times 10^{-4}$		Adequate (greater than 2.5 m/sec)	Aspirations Psych. Preussischen Met. Ins., 1930.
$6.56 \times 10^{-4}$	$5.79 \times 10^{-4}$	Idem	Jelineks Psych. Jafeln, 1929.
$12.0 \times 10^{-4}$	$10.6 \times 10^{-4}$	Inadequate, Almost 0 m/sec	Jelineks Psych. Jafeln, 1929.
$8.00 \times 10^{-4}$	$7.06 \times 10^{-4}$	Inadequate, 1 to 1.5 m/sec	Jelineks Psych. Jafeln, 1929.
$6.50 \times 10^{-4}$	$5.80 \times 10^{-4}$	Idem	Meteorologie Nationale France, Nouvelles Tables, (1956).

If one attempts to evaluate  $T_w$  by substituting Eq. 1 into Eq. 2 the problem is that  $T_w$  in the first member is affected by an exponential function and in the second by a linear function. So it is not possible to apply conventional algebraic methods in order to write  $T_w$  as a function of  $T$  and  $RH$ . This problem arises when explicit values of  $T_w$  are not given in the climatological records but  $T$  and  $RH$  values are readily available.

## 2. Vapor Pressure Evaluation

Equation 1 is not an exact expression for  $e_s$ . In comparison with experimental values reported by Byers (1959, p.158) the root mean squared error (RMSE) of Eq. 1 is 0.8 hPa between temperatures of  $-10^{\circ}\text{C}$  to  $50^{\circ}\text{C}$ .

Here it is proposed the following polynomial expression:

$$e_s = aT^3 + bT^2 + cT + d \quad (3)$$

where  $a = 6.6 \times 10^{-4}$ ,  $b = 4.6 \times 10^{-3}$ ,  $c = 4.58 \times 10^{-1}$ ,  $d = 6.63$ , and  $e_s$  units are hPa.

Other authors have proposed empirical or semi-empirical methods to evaluate  $e_s$ . Table 2 shows comparisons by means of the Pearson linear correlation coefficient ( $R$ ) and the RMSE between experimental values of  $e_s$  (Byers, 1959) and some physical and empirical methods for each degree of temperature.

It is important to note that a non-exponential expression for  $e_s$  is an aid in the evaluation of  $T_w$  as a function of  $T$  and  $RH$  from Eq. 2. So, in view of the goodness of the polynomial expression here proposed (Eq. 3), it is useful to find the function ( $f$ ) for  $T_w = f(T, RH, P)$ .

Table 2. Comparison between the experimental values of  $e_s$  (Byers, 1959; p. 158), the Clausius - Clapeyron values and some empirical methods, from  $-10$  to  $50^{\circ}\text{C}$ .

Author	Range ( $^{\circ}\text{C}$ )	Pearson linear correl. coef.	RMSE (hPa)
Clausius-Clapeyron <sup>1</sup>	-10 to 50	0.9999	0.80
Adem, 1967 <sup>2</sup>	-10 to 50	0.9998	0.80
Lowry and Lowry, 1989 <sup>3</sup>	5 to 35	0.9996	0.70
Rosenberg, 1974 <sup>4</sup>	-10 to 50	0.9982	1.00
Steadman, 1979 <sup>5</sup>	-10 to 50	0.9997	0.80
Steadman, 1979	-10 to 40	0.9995	0.70
Equation 3	-10 to 50	0.9997	0.60

1) Equation 1.

2)  $e_s = 6.115 + 0.42915 T + 1.4206 \times 10^{-2} T^2 + 3.046 \times 10^{-4} T^3 + 3.2 \times 10^{-6} T^4$ .

3)  $e_s = 8.51 + 0.037 T^2$ .

4)  $\log_{10} e_s = 0.82488 + 0.02604 T$ .

5)  $e_s = 6.46 + 0.555 T + 7.1 \times 10^{-4} T^3$

### 3. The Evaluation of $T_w$

#### 3.1 Polynomial expression

By using Eq. 3 to calculate  $e_{sw}$  and Eq. 1 to calculate  $e_s$ , the psychrometric equation (Eq. 2) can be written:

$$T_w^3 + \frac{b}{a}T_w^2 + \frac{c + kp}{a}T_w + \frac{d - (RH/100)e_s - kpT}{a} = 0 \quad (4)$$

By substitution of:

$$T_w = tw - \frac{b}{3a},$$

$$S = \frac{c + kp}{a} - \frac{1}{3} \left( \frac{b}{a} \right)^2$$

$$Q = \frac{2b^3}{27a^3} - \frac{b(c + kp)}{3a^3} + \frac{d - (RH/100)e_s - kpT}{a}$$

Eq. 4 can be rewritten so:

$$tw^3 + Stw + Q = 0 \quad (5)$$

The discriminant  $[(Q/2)^2 + (S/3)^3] > 0$  for atmospheric values near the ground. According the Cardano's method (Gellert *et al.*, 1975; p. 97-99) Eq. 5 has one real solution and two conjugate complex solutions. The real solution is:

$$T_w = \left\{ -\frac{Q}{2} + \left[ \frac{Q^2}{4} + \frac{S^3}{27} \right]^{1/2} \right\}^{1/3} + \left\{ -\frac{Q}{2} - \left[ \frac{Q^2}{4} + \frac{S^3}{27} \right]^{1/2} \right\}^{1/3} + \frac{b}{3a} \quad (6)$$

#### 3.2 Iterative methods

In order to establish comparisons with the estimation made by means of the process of section 3.1, a computation program to evaluate  $T_w$  from Eq. 2 was implemented as follows:

$$T_w = T - \frac{e_{sw_0} - (RH/100)e_s}{kp}, \quad (7)$$

where  $T$  is the dry-bulb temperature in °C,  $RH$  is the relative humidity in %,  $e_s$  is the saturation vapor pressure (in hPa) found through Eq. 1,  $k = 6.53 \times 10^{-4} \text{ C}^{-1}$  and  $P$  is the atmospheric pressure in hPa.  $e_{sw_0}$  is the obtained value of Eq. 1 applied to an initial proposed value of  $T_{w_0} = T - 5 \text{ °C}$ , which can be approximated to correct  $T_w$  thus:

a) If  $|T_w - T_{w_0}| \leq 0.1 \text{ °C}$ , then the evaluation of  $T_w$  is "correct";

b) If  $(T_w - T_{w_0}) > 0.1 \text{ °C}$ , then a new value for  $T_{w_0}$ , i.e.  $T_{w_0} = T - 5 \text{ °C} + N \times 10^{-3}$ , is proposed where  $N$  is a whole number that indicates the order of iterative recurrence needed to obtain the condition (a).

c) If  $(T_w - T_{w_0}) < -0.1$  °C, then  $T_{w_0} = T - 5$  °C  $- N \times 10^{-3}$  and the iteration runs until (a) is obtained.

#### 4. Comparison Methods

The results of the application of the method showed in section 3.1 were compared and corrected by means of regular linear regression with the values of iterative recurrences of section 3.2, for pressure levels of 1000, 950, 900, 850, 800, 750, 700 and 650 hPa, dry-bulb temperatures from -10 to +50°C for ranges of 1°C, and RH = 10, 20, 30,...90%.

Thus by substitution of constants and by the use of the above mentioned linear regression correlations the numerical expressions used in Eq. 6 were:

$$S = 662.23 + 0.97p \quad (8a)$$

$$Q = 8264.65 - 1480.45(RH/100)e_s - 0.966pT \quad (8b)$$

$$\frac{b}{3a} = -1^\circ C \quad (8c)$$

For the 4392 points of comparison the correlation coefficient was of the order of 0.999 and the RMSE was 0.7 °C. Table 3 shows the goodness of fit as a function of the pressure levels.

Table 3. Goodness of fit of the polynomial method (Eqs. 6 and 8) for the evaluation of  $T_w$  at various pressure levels.

Pressure level (hPa)	Pearson linear Correl. Coef. (R)	Coef. of Determination (R <sup>2</sup> )	RMSE (°C)
1000	0.9996	0.9991	0.4
950	0.9995	0.9990	0.5
900	0.9994	0.9988	0.5
850	0.9993	0.9985	0.6
800	0.9990	0.9980	0.6
750	0.9987	0.9974	0.7
700	0.9982	0.9966	0.9
650	0.9978	0.9956	1.0

## 5. Conclusions

Eq. 3 is a good expression for the evaluation of  $e_s$  for water-vapor phases between  $-10^{\circ}\text{C}$  to  $50^{\circ}\text{C}$ . In view of the fact that this equation is not an exponential expression, it can be used to obtain  $Tw$  as a function of  $T$ ,  $RH$  and  $p$  by means of the equations 6 and 8, with high goodness of fit, mainly at low pressure levels (below 850 hPa), where the RMSE is almost  $0.5^{\circ}\text{C}$ . The further advantage is that these procedures may be implemented with a pocket calculator.

## Aknowledgements

The author wishes to tank Mr. P. Mosiño and Dr. E. Jáuregui for their suggestions, and Mrs. L. Calatayud for her help with translation. Thanks are also due to Mr. J. Ruiz-Barradas, who did the computer work. This paper was supported by National Council of Science and Technology (CONACYT) and the Atmospheric Sciences Center (UNAM) under project O521T.

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